

ON SOME SYMMETRY AXIOMS IN RELATIVITY THEORIES

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Abstract: *In this paper we review two symmetry axioms of special relativity and their connections to each other together with their role in some famous predictions of relativity theory, such as time dilation, length contraction, and the twin paradox. We also discuss briefly counterparts of these symmetry axioms in general relativity and formulate a conjecture, namely that without them the axioms of general relativity would capture general relativistic spacetimes only up to conformal equivalence.*

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1. INTRODUCTION

Special relativity is axiomatic in its spirit from the beginning since Einstein introduced special relativity as a theory of two informal postulates in his famous 1905 paper. One of these postulates was Galileo's principle of relativity (generalized to all physical phenomena), implying that there is no physical phenomenon that distinguishes some inertial observers; the other was the natural assumption that the speed of light is the same in every direction at least in one inertial frame of reference, see Einstein (1905, Sec. 2).

Even in the non-axiomatic approaches to general relativity, a spacetime is given as a triple of a manifold M , a Lorentzian metric g , and an energy-momentum tensor T satisfying Einstein's Field Equations and often some extra conditions, such as the causality or the energy conditions. So general relativity is also axiomatic in its spirit because in this context it is natural to think of M , g , and T as basic concepts and of their basic properties (e.g., M is being a manifold) together with Einstein's Field Equations and the extra conditions as axioms.

Formal axiomatizations of relativity theories (both special and general) also have an extended literature, see, e.g., Ax (1978), Benda (2008), Benda (2015), Goldblatt (1987), Guts (1982), Mundy (1986), Pambuccian (2007), Schutz (1981), Szabó (2009) to mention only a few. One of these formal approaches is the one developed by the research team/school of Hajnal Andréka and István Németi. Here we will stay within the main framework of this team. However, comparing the axiomatizations of different frameworks is also an interesting and important research direction, see, e.g., Andréka and Németi (2014), Barrett and Halvorson (2015), Weatherall (2014), Rosenstock et. al. (2015).

In this paper, after recalling the main axiomatic framework of the Andréka–Németi school, we review two of the numerous symmetry axioms appearing in this approach. We also review how these symmetry axioms are related to each other and their role in predictions of relativity theory, such as time dilation, length contraction, and the twin paradox. Finally, we discuss the counterparts of these symmetry axioms in general relativity and formulate a conjecture about what happens if we leave out these counterparts from the axiom system of general relativity, see Conjecture 4 on p. 12.

Apart from slight strengthening of some theorems resulting from a natural generalization of axiom system *SpecRel*₀ and the simple observations of Propositions 1, 2 and 3, all the theorems and propositions of this paper can be found scattered in the cited references. The main aim of this paper is to investigate the roles and intuitive meanings of the two most often used symmetry axioms of the Andréka–Németi school in a way that can easily be understood even by nonspecialists.

For methodological reasons, we work in a formal axiomatic framework, which among others is beneficial because it forces us to formulate unambiguous basic assumptions with a clear meaning. See Andréka et. al. (2002, pp. 1245-1252) and Székely (2009, Sec. 11.) for more methodological details why first-order logic is an ideal logic for developing axiomatic frameworks for relativity theories. A practical advantage of using first-order logic is that the theorems can be machine verified, see, e.g., Govindarajulu et. al. (2015), and Stannett and Németi (2014). See Friend (2015) for epistemological significance of the Andréka–Németi approach.

To make our formulas easy to read even to non-logicians, we try to use only simple and natural notations, e.g., the logical connective “implies” is denoted by \implies , and logical connective “and” is denoted by comma. Quantifiers “for all” and “exists” are denoted by the usual symbols \forall and \exists , respectively.

2. AXIOMATIC FRAMEWORK

2.1. Basic concepts. First we fix a set of basic concepts about which we will formulate some basic assumptions (axioms).¹ Here we use the main framework of the Andr ka–N meti school. That is, we have two sorts of basic concepts: *Bodies* (things that move) B and *Quantities* (numbers that are used to coordinatize the moving bodies) Q . We have two kinds of distinguished bodies *inertial observers* (or inertial coordinate systems) IOb and *light signals* (or photons) Ph . To put an algebraic structure on the quantities, we take the usual operations $+$, \cdot and the ordering $<$ as basic concepts. Finally, we connect the physical sort (of bodies) and the mathematical sort (of quantities) by the worldview relation W .

We use the worldview relation W to express how observers associate coordinates to events (i.e., meetings of bodies). This is done by translating basic relation $W(o, b, x, y, z, t)$ to natural language as “Observer o coordinatizes body b at space location (x, y, z) at instant t .”

The two key concepts events and worldlines of relativity theory can easily be defined from W as follows. The *event* coordinatized by observer o at spacetime location $\bar{x} = (x, y, z, t)$ is the set of those bodies that are coordinatized at \bar{x} by o , i.e.,

$$ev_o(\bar{x}) \stackrel{\text{def.}}{=} \{b : W(o, b, \bar{x})\}.$$

The *worldline* of body b according to observer o is the set of coordinate points where b is coordinatized by o , i.e.,

$$wline_o(b) \stackrel{\text{def.}}{=} \{\bar{x} : W(o, b, \bar{x})\}.$$

For any two coordinate points $\bar{x}, \bar{y} \in Q^4$, let us use the following notations for the *spatial distance* and the *time difference*:

$$dist(\bar{x}, \bar{y}) \stackrel{\text{def.}}{=} \sqrt{(x_1 - y_1)^2 + (x_2 - y_2)^2 + (x_3 - y_3)^2} \text{ and } time(\bar{x}, \bar{y}) \stackrel{\text{def.}}{=} |x_4 - y_4|. \text{ }^2$$

The *speed* of body b according to observer i can be defined as follows,

$$v_i(b) = v \stackrel{\text{def.}}{\iff} (\exists \bar{x}, \bar{y} \in wline_i(b)) [space_i(\bar{x}, \bar{y}) = v \cdot time(\bar{x}, \bar{y})]$$

if $wline_i(b)$ is a subset of a line, e.g., if b is a light signal, i is an inertial observer, and axioms $AxPh$ and $AxEField$ (below) are assumed.

¹We do not fix the basic concepts once and for all. We fix them just for creating a framework to formulate certain axioms and axiom systems. The axiomatic method also has the flexibility of modifying the basic concepts. Moreover, comparing axiom systems formulated using different basic concepts is an interesting research area, see, e.g., Andr ka and N meti (2014).

²By axiom $AxEField$ (see below) we have a strong enough algebraic structure on the quantities to define subtraction and square root.

The *worldview transformation* between observer o and o' connects coordinate points \bar{x} and \bar{x}' iff the event coordinatized by o at \bar{x} is the same as the one coordinatized by o' at \bar{x}' , i.e.,

$$w_{oo'}(\bar{x}) = \bar{x}' \stackrel{\text{def.}}{\iff} ev_o(\bar{x}) = ev_{o'}(\bar{x}').$$

By the above definition worldview transformations are only binary relations but after assuming some axioms they become transformations, see, e.g., Theorem 5.

2.2. Axioms for special relativity. Einstein's two original postulates immediately imply that the speed of light is the same for every inertial observer because this is true for one of the inertial observers by the second postulate, and there is no distinguished inertial observer by the first postulate. This property of light signals is basically the only nontrivial assumption we need to capture the kinematics of special relativity.

AxPh: For every inertial observer i , there is a finite speed c_i such that all light signals move with speed c_i according to i , and it is possible to send out a light signal with this speed c_i in every direction everywhere, i.e.,³

$$(\forall i \in IOb)(\exists c_i \in Q)(\forall \bar{x}, \bar{y} \in Q^4) \left[(\exists p \in Ph) [\bar{x}, \bar{y} \in wline_i(p)] \iff dist(\bar{x}, \bar{y}) = c_i \cdot time(\bar{x}, \bar{y}) \right].^4$$

Not just to prove their usual properties, but even to be able to define the concepts of spatial distance and time difference, we need some assumptions about the quantities. Therefore, we assume some algebraic properties of real numbers.

AxEField: $(Q, +, \cdot, <)$ is an Euclidean field.⁵

Also *AxPh* has its intended meaning only if we assume that inertial observers coordinatize the same outside reality. Therefore, we also need the following assumption.

AxEv: Inertial observers coordinatize the same events, i.e.,

$$(\forall i, i' \in IOb)(\forall \bar{x} \in Q^4)(\exists \bar{x}' \in Q^4) [ev_i(\bar{x}) = ev_{i'}(\bar{x}')].$$

Finally, to make it easier speaking about the motion of coordinate systems in the usual way by referring to the image of their time axes in other coordinate systems,

³It is more natural to assume that the value c is the same for all inertial observers in a separate axiom, see axioms $Ax(c_i = c_j)$ and $Ax(c = 1)$ on p. 8.

⁴This formula literally says that for any inertial observer i there is a quantity c_i such that a potential light signal p can be located in coordinate points \bar{x} and \bar{y} according to i if and only if the corresponding speed is c_i .

⁵That is $(Q, +, \cdot)$ is a field, $(Q, <)$ is a linearly ordered set, $x < y \implies x + z < y + z$ and $0 < x, 0 < y \implies 0 < x \cdot y$, and every positive number has a square root, i.e., $(\forall x > 0) \exists y [x = y^2]$.

we assume that the worldline of every inertial observer is the time axis in her/his own coordinate system.

AxSelf: Inertial observers are stationary in the origin of their own coordinate systems, i.e.,

$$(\forall i \in IOb) \forall xyz t [(x, y, z, t) \in wline_i(i) \iff x = y = z = 0].$$

The four simple axioms above are enough to capture special relativity in a qualitative way. So let us introduce *SpecRel*₀ as the axiom system containing the four axioms above:

$$SpecRel_0 \stackrel{\text{def.}}{=} AxPh + AxEFIELD + AxEv + AxSelf.$$

2.3. Consequences of *SpecRel*₀. The first thing the reader might spot is that there is no axiom in *SpecRel*₀ stating that inertial observers move uniformly. This is so because this statement follows from the axioms of *SpecRel*₀:

Theorem 1.⁶

$$\begin{aligned} SpecRel_0 \implies (\forall i, j \in IOb) (\forall \bar{x}, \bar{y}, \bar{z} \in wline_i(j)) [\bar{x} \neq \bar{y} \\ \implies \exists \lambda [\lambda \cdot (\bar{x} - \bar{y}) = \bar{z} - \bar{y}]] \end{aligned}$$

We made no explicit restriction on the speed of the inertial observers, but *SpecRel*₀ implies that they cannot move faster than the speed of light with respect to one another.

Theorem 2. *SpecRel*₀ $\implies (\forall i, j \in IOb) (\forall p \in Ph) [v_i(j) < v_i(p)]$

A direct proof of Theorem 2 from slightly stronger axioms can be found, for example, in Andr  ka et. al. (2012a).

By Theorem 3 below *SpecRel*₀ contradicts the Newtonian notion of absolute time and replaces it with an observer dependent one. For a formal statement and a direct proof from slightly stronger axioms, see, e.g., Andr  ka et. al. (2007, Thm. 11.4, pp. 626-630).

Theorem 3. *SpecRel*₀ \implies “For any two relatively moving inertial observers, there are pairs of events which are simultaneous for one of the observers but not for the other.”

Time dilation and length contraction are two famous predictions of special relativity. Even though *SpecRel*₀ does not imply the exact rates of the dilation and the contraction, it predicts these effects qualitatively in the following way.

⁶By *AxSelf* this theorem is a direct consequence of Theorem 5 below.

Theorem 4.

- (1) $SpecRel_0 \implies$ “At least one of two relatively moving inertial observers sees that the other’s clocks are slowed down.”
- (2) $SpecRel_0 \implies$ “At least one of two relatively moving inertial spaceships shrinks according to the other.”

See Andr  ka et. al. (2002, Sec. 2.5) for precise formulations of these statements.

All the theorems above can also be derived from Theorem 5 below. To state this theorem let us recall that, a map $P : Q^4 \rightarrow Q^4$ is a *Poincar   transformation* if it is an affine bijection with the following property

$$time(\bar{x}, \bar{y})^2 - dist(\bar{x}, \bar{y})^2 = time(\bar{x}', \bar{y}')^2 - dist(\bar{x}', \bar{y}')^2$$

for all $\bar{x}, \bar{y}, \bar{x}', \bar{y}' \in Q^4$ for which $P(\bar{x}) = \bar{x}'$ and $P(\bar{y}) = \bar{y}'$.

Since Poincar   transformations are the standard coordinate transformations in special relativity, Theorem 5 below basically says that $SpecRel_0$ implies these standard coordinate transformations up to changing units of measurement.⁷

Theorem 5. $SpecRel_0 \implies$ “The worldview transformations between inertial observers are Poincar   transformations up to changing the units of measurement and permuting all coordinates by a field automorphism.”⁸

Theorem 5 follows immediately from Andr  ka et. al. (2012b, Theorem 7.8) by (Ibid Proposition 7.19).

3. SYMMETRY AXIOMS IN SPECIAL RELATIVITY

A great many symmetry axioms were formulated and investigated in Andr  ka et. al. (2002, Sec. 2.8 and Sec. 3.9). Here we highlight only two of them, both of which are about harmonizing the units of measurement of different inertial observers.

So how can we test if two different inertial observers use the same units of measurement or not?

⁷Permuting the coordinates by a field automorphism can also be considered as an unusual way of changing the units of measurement if the underlying field has a nontrivial automorphism at all. For example, the field of rational numbers or the field of real numbers does not have a nontrivial automorphism.

⁸That is, the coordinate transformation between inertial observers i and i' becomes a Poincar   transformation after multiplying all the space coordinates by a positive number, the time coordinates with a possibly different number, and maybe also transforming the coordinate system of i by a mapping (x, y, z, t) to $(\phi(x), \phi(y), \phi(z), \phi(t))$, where ϕ is an automorphism of $(Q, +, \cdot)$ and doing the same with the coordinates of i' but with possibly different numbers and automorphism.

3.1. Symmetry of space. To compare the units of two different inertial observers measuring spatial distances, we can ask them to determine the distances of events which are simultaneous for both of them. If they get the same distances for these events, then they use the same units. Otherwise, the one who uses the smaller units gets bigger numbers for the distances. In the next axiom, we can use this idea to formulate that different inertial observers use the same units to measure spatial distances.

AxSymDist: Inertial observers agree as for the spatial distance between events if these events are simultaneous for both of them, i.e.,

$$(\forall i, i' \in IOb)(\forall \bar{x}, \bar{y}, \bar{x}', \bar{y}' \in Q^4) \left[ev_i(\bar{x}) = ev_{i'}(\bar{x}'), ev_i(\bar{y}) = ev_{i'}(\bar{y}'), \right. \\ \left. time(\bar{x}, \bar{y}) = 0, time(\bar{x}', \bar{y}') = 0 \implies dist(\bar{x}, \bar{y}) = dist(\bar{x}', \bar{y}') \right].$$

3.2. Symmetry of time. Another idea to compare the units of measurement of different inertial observers is to ask them to compare the ticking rate of each others' clocks. If they see each others' clock behaving the same way, e.g., slowed down with the same rate, then they use the same units to measure time differences. Otherwise, the one who uses the smaller unit considers the other's clock slowed down more (he would consider the other's clocks slow even if they are stationary with respect to each other). The next axiom uses this idea to formulate that different inertial observers use the same units to measure time differences.

AxSymTime: Any two inertial observers see each others' clocks slowed down with the same rate, i.e.,

$$(\forall i, i' \in IOb) \forall t \left[time(w_{ii'}(t \cdot \bar{1}_t), w_{ii'}(\bar{o})) = time(w_{i'i}(t \cdot \bar{1}_t), w_{i'i}(\bar{o})) \right],$$

where $\bar{o} = (0, 0, 0, 0)$ and $\bar{1}_t = (0, 0, 0, 1)$.

3.3. Connections between the symmetry axioms. The value of the speed of light depends on the units of measurement inertial observers choose to measure space and time. After having *AxPh* assumed, assuming that the speed of light is 1 only means that inertial observers measure spatial distance in the units corresponding to their time unit. For example, if somebody measures time in years he only has to measure distance in light-years to have 1 for the speed of light, etc.

If *AxPh* and *AxEField* are assumed, then we can ensure that the speed of light is the same for different inertial observers by the following axiom.

$Ax(c_i = c_j)$: Inertial observers use units to measure time differences and spatial distances such that the speed of light is the same for all of them, i.e.,

$$(\forall i, j \in IOb)(\forall p \in Ph)[v_i(p) = v_j(p)].$$

Likewise, if $AxPh$ and $AxEField$ are assumed, then we can set the speed of light to 1 for all the inertial observers by the following axiom.

$Ax(c = 1)$: Inertial observers use units to measure time differences and spatial distances such that the speed of light is one for all of them, i.e.,

$$(\forall i \in IOb)(\forall p \in Ph)[v_i(p) = 1].$$

Originally, $Ax(c = 1)$ was part of $AxPh$, see, e.g., Andr ka et. al. (2002, AxE on p. 51) or Andr ka et. al. (2007, p. 621). Later it was moved to the corresponding symmetry axiom because it is strongly related to the units of measurement chosen by inertial observers, see, e.g., Andr ka et. al. (2012a). Here we have introduced $Ax(c = 1)$ as an axiom in its own right because it is easier to understand the roles and intuitive meanings of $AxSymDist$ and $Ax(c = 1)$ if they are separated.

Let us introduce axiom system $SpecRel$ as follows:

$$SpecRel \stackrel{\text{def}}{=} SpecRel_0 + AxSymDist + Ax(c = 1).$$

Theorem 6. *$SpecRel$ is a consistent axiom system. Moreover, there is a model of $SpecRel$ where for every inertial observer i and Poincar  transformation P there is an inertial observer j such that $w_{ij} = P$.*

See Andr ka et. al. (2002, Sec. 3.6) and Andr ka et. al. (2007, Sec. 2.5) for constructions proving theorem 6.

By Proposition 1 below, $AxSymTime$ and $AxSymDist$ are independent from the rest of the axioms of $SpecRel$ and they are not equivalent if only $SpecRel_0$ is assumed.

Proposition 1.

- (1) $SpecRel_0 + Ax(c = 1) \not\Rightarrow AxSymTime$
- (2) $SpecRel_0 + Ax(c = 1) \not\Rightarrow AxSymDist$
- (3) $SpecRel_0 + AxSymTime \not\Rightarrow AxSymDist$
- (4) $SpecRel_0 + AxSymDist \not\Rightarrow AxSymTime$

On the proof. All of the items of this proposition are proved by constructing appropriate models. We can start from any model of $SpecRel$ containing enough (at least two) observers. For example, the one whose existence is stated by Theorem 6.

To prove items 1. and 2., we should construct a model of $SpecRel_0 + Ax(c = 1)$, where $AxSymTime$ and $AxSymDist$ do not hold. If we modify an inertial observer's coordinate system in a model of $SpecRel_0 + Ax(c = 1)$ by multiplying all the coordinates with the same positive number, we get another model of $SpecRel_0 + Ax(c = 1)$. Using this modification it is easy to make $AxSymTime$ and $AxSymDist$ invalid in the models of $SpecRel$ without changing the validity of the rest of the axioms.

To prove items 3. and 4., we should make one of $AxSymTime$ and $AxSymDist$ invalid without changing the validity of the other and the axioms of $SpecRel_0$. This can be done easily by scaling only the time or only the space coordinates of the inertial observers.⁹ \square

By Proposition 2 below, $AxSymTime$ and $AxSymDist$ are equivalent assuming $Ax(c_i = c_j)$ and $SpecRel_0$.

Proposition 2. $SpecRel_0 + Ax(c_i = c_j) \implies (AxSymTime \iff AxSymDist)$

On the proof. If $Ax(c = 1)$ is also satisfied, both $AxSymTime$ and $AxSymDist$ are equivalent to the statement that the worldview transformations are Poincaré transformations. The other cases can easily be reduced to the $c = 1$ case by multiplying all the observer's time coordinates by factor c (i.e., the speed of light that is the same for all inertial observers by axiom $Ax(c_i = c_j)$). \square

By Proposition 3 below, $AxSymTime$ and $AxSymDist$ imply that different inertial observers agree on the exact value of the speed of light if $SpecRel_0$ is assumed.

Proposition 3. $SpecRel_0 + AxSymTime + AxSymDist \implies Ax(c_i = c_j)$

On the proof. Using Theorem 5, it is easy to prove this proposition. Axioms $AxSymTime$ and $AxSymDist$ fix both the time and space units of measurement of inertial observers. Without changing these units what remains from the possible worldview transformations characterized by Theorem 5 leaves no flexibility to have inertial observers with different values for the speed of light. \square

3.4. The role of the symmetry axioms in certain predictions of relativity. By Theorem 7 below, $SpecRel$ captures time dilation and length contraction predictions of special relativity even quantitatively.

⁹Of course, this construction also makes axiom $Ax(c = 1)$ invalid by Proposition 2.

Theorem 7.

- (1) *SpecRel* \implies “Relatively moving inertial observers see that each others’ clocks are slowed down exactly by the Lorentzian contraction factor.”
- (2) *SpecRel* \implies “Relatively moving inertial observers see each others’ spatial distances shrinking exactly by the Lorentzian contraction factor in the direction of motion.”

See Andr ka et. al. (2007, Thm. 11.6, pp. 631-635) for precise formalization of these statements and direct proofs from the axioms.

The so called twin paradox is the prediction of relativity stating that between two distinct meeting points inertial observers always measure more time than non-inertial ones do. Since it considers only inertial observers, *SpecRel* is not strong enough to investigate the twin paradox in this form.¹⁰ However, we can simulate the accelerated twin by two inertial observers: a leaving one and a returning one. With this standard trick we can introduce an inertial version of the twin paradox where in the simulated twin paradox situations (cf. Figure 1) the stay-at-home inertial twin always measures more time than his leaving and returning inertial sisters together.

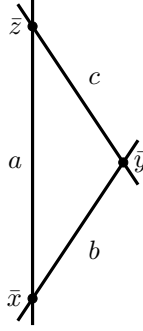


FIGURE 1. Illustration for the formula *TwP*

TwP: In the situations depicted by Figure 1 inertial observer *a* measures more time between the events at \bar{x} and \bar{y} than inertial observers *b* and *c* together, i.e.,

$$\begin{aligned}
 &(\forall a, b, c \in IOb)(\forall \bar{x}, \bar{z} \in wline_a(a))(\forall \bar{y} \notin wline_a(a)) \\
 &\quad \left[\bar{x} \in wline_a(b), \bar{y} \in wline_a(b) \cap wline_a(c), \bar{z} \in wline_a(c) \right. \\
 &\quad \left. \implies time(\bar{x}, \bar{z}) > time(w_{ab}(\bar{x}), w_{ab}(\bar{y})) + time(w_{ac}(\bar{y}), w_{ac}(\bar{z})) \right]
 \end{aligned}$$

¹⁰See Madar sz et. al. (2006) and Sz kely (2009, Sec. 7) for an axiomatic investigation of the twin paradox in this form within a theory of accelerated observers.

By Theorem 8 below, $SpecRel$ implies TwP but not without the symmetry axioms.

Theorem 8.

- (1) $SpecRel_0 + Ax(c = 1) \not\Rightarrow TwP$
- (2) $SpecRel_0 + Ax(c = 1) \Rightarrow (AxSymTime \Rightarrow TwP)$

In view of Theorem 8, whether TwP also implies $AxSymTime$ or not is a natural question asked by Andr  ka et. al. (2002, Question 4.2.17). By the intuitive insight of this paper about the role of axiom $AxSymTime$, it is not surprising that TwP cannot fulfill that role and hence it does not imply $AxSymTime$.

Proposition 4. $SpecRel_0 + Ax(c = 1) + TwP \not\Rightarrow AxSymTime$

See Sz  kely (2010) and Sz  kely (2009 Sec. 4) for proof of Theorem 8 and Proposition 4 and a detailed investigation of TwP , e.g., a geometrical characterization of the models of $SpecRel_0 + Ax(c = 1)$ where TwP holds.

4. SYMMETRY AXIOMS IN GENERAL RELATIVITY

Both in the special and the general theories of relativity, this kind of axiomatic approach to symmetries based on comparing the units of measurement of different observers is original to the Andr  ka–N  meti team. To discuss the general relativistic counterparts of symmetry axioms $AxSymTime$ and $AxSymDist$, let us see first how general relativity emerges in the axiomatic framework of the team.

The transition from special relativity to general relativity in terms of axioms basically goes by assuming the localized versions of the same axioms, see, e.g., Andr  ka et. al. (2012a). The corresponding axiom system is called $GenRel$. Theorem 9 below says that the axiom system $GenRel$ captures general relativity well. For more details, e.g., precise statement, proof, and refinements, see, e.g., Andr  ka et. al. (2013).

Theorem 9 (Completeness Theorem). *$GenRel$ is complete with respect to the standard models of general relativity, i.e., to Lorentzian manifolds.*

How are the axioms of $GenRel$ related to the axioms of $SpecRel$?

For every axiom Ax of $SpecRel$, there is an axiom Ax^- in $GenRel$, such that Ax^- captures the same idea locally for arbitrary observers as Ax does globally for inertial ones. The only exception is axiom $AxDiff$ of $GenRel$ because it localizes not

an axiom but the theorem of *SpecRel* stating that the worldview transformations between inertial observers are affine transformations.

In *GenRel*, the localized version of axiom $Ax(c = 1)$ is included in $AxPh^-$. Both $AxSymDist$ and $AxSymTime$ can be localized, but the localization of $AxSymTime$ is simpler, see Székely (2009, pp. 96-97). $AxSymTime^-$ intuitively says the following.

$AxSymTime^-$: Any two observers see each others' clocks behaving in the same way at an event of meeting.

The standard approach to symmetries in general relativity is based on local diffeomorphisms preserving some geometrical notions, such as the metric or geodesics, see, e.g., Hall (2004, Sec. 10). Even though it is a natural open question how these standard concepts of symmetries can be captured within the first-order logic framework of the Andr  ka-N  meti team, here we concentrate on another interesting question, namely on what role may the symmetry axioms play in *GenRel*.

In *SpecRel* the only flexibility we get by leaving out the symmetry axioms is that different inertial observers may use different units of measurements, see Theorem 5. Therefore, it is natural to conjecture that in *GenRel* the localized version of the symmetry axioms will have similar roles, i.e., without them the rest of the axioms will still capture the standard models of general relativity up to conformal equivalence (i.e., up to changing of units of measurements locally).¹¹

Conjecture. *If we leave out $AxSymTime^-$ from *GenRel* it will still capture Lorentzian manifolds up to conformal equivalence.*

That is, even though no unique Lorentzian metric can be defined without axiom $AxSymTime^-$, the manifold of events and the tangent space is definable the same way as in Andr  ka et. al. (2013) and probably there will be some definable geometric object in the tangent space that can capture the conformal equivalence classes of Lorentzian manifolds.

5. CONCLUDING REMARKS

We have seen that symmetry axioms $AxSymTime$ and $AxSymDist$ ensure that different inertial observers use the same units of measurement to measure time and

¹¹Since $AxPh^-$ contains the localized version of $Ax(c = 1)$, the units for measuring time and space remain intertwined even if we omit $AxSymTime^-$. The nonstandard unit changing transformations induced by field automorphisms in Theorem 5 are conjectured to be ruled out by $AxDiff$.

space, respectively. This insight it helps to understand why these symmetry axioms are needed to get the exact Lorentzian contraction factors and why $SpecRel_0$ implies time dilation and length contraction only qualitatively. Also after understanding the exact roles of these symmetry axioms, it becomes clearer why they are needed to prove the twin paradox (TwP), and why TwP is not strong enough to imply them.

It is natural to conjecture that the localized version of these symmetry axioms will have a similar role in general relativity, but to check this conjecture requires further investigations.

6. REFERENCES

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